

WIDE RADIO BEAMS FROM γ -RAY PULSARS

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 Received 2010 February 24; accepted 2010 May 11; published 2010 May 24

ABSTRACT

We investigate the radio and γ -ray beaming properties of normal and millisecond pulsars (MSPs) by selecting two samples from the known populations. The first, Sample G, contains pulsars which are detectable in blind searches of γ -ray data from the *Fermi* Large Area Telescope. The second, Sample R, contains pulsars detectable in blind radio searches which have spin-down luminosities $\dot{E} > 10^{34}$ erg s $^{-1}$. We analyze the fraction of the γ -ray-selected Sample G which have detectable radio pulses and the fraction of the radio-selected Sample R which have detectable γ -ray pulses. Twenty of our 35 Sample G pulsars have already observed radio pulses. This rules out low-altitude polar-cap beaming models if, as is currently believed, γ -ray beams are generated in the outer magnetosphere and are very wide. We further find that, for the highest- \dot{E} pulsars, the radio and γ -ray beams have comparable beaming factors, i.e., the beams cover similar regions of the sky as the star rotates. For lower- \dot{E} γ -ray emitting pulsars, the radio beams have about half of the γ -ray sky coverage. These results suggest that, for high- \dot{E} young and MSPs, the radio emission originates in wide beams from regions high in the pulsar magnetosphere, probably close to the null-charge surface and to the γ -ray emitting regions. Furthermore, it suggests that for these high- \dot{E} pulsars, as in the γ -ray case, features in the radio profile represent caustics in the emission beam pattern.

Key words: pulsars: general

1. INTRODUCTION

The observed pulses from pulsars are interpreted as an emission beam, fixed relative to the rotating neutron star, sweeping across the Earth. The magnetospheric emission-region structure, the viewing angle with respect to the rotation axis, and the angle between the magnetic and rotation axes define the shape of the observed pulse profile. Pulsar radio emission is generally thought to originate within a few percent of the light-cylinder radius (R_{LC}) above the polar caps (e.g., Lyne & Manchester 1988). Many pulsars exhibit a “radius-to-frequency” mapping (Cordes 1978) which is indicative of polar cap emission. The increasing pulse width at lower radio frequencies is interpreted as emission at increasing heights above the poles along more widely directed field lines. The radio profiles of many young pulsars and millisecond pulsars (MSPs) have an “interpulse,” a second pulse component separated from the main pulse by close to 180° of the pulse phase, or widely separated pulse components. In most cases, such pulsars do not exhibit radius-to-frequency mapping, that is, they have frequency-independent pulse component separations (Hankins & Fowler 1986).

The known population of γ -ray pulsars has recently been greatly expanded by the 6 month data release of the *Fermi* Large Area Telescope (LAT; Abdo et al. 2010b, hereafter A10). This sample of 46 γ -ray pulsars includes both young pulsars and MSPs, all with high spin-down luminosities:

$$\dot{E} = \frac{4\pi^2 I \dot{P}}{P^3} > 10^{33} \text{ erg s}^{-1}, \quad (1)$$

where $I = 10^{45}$ gm cm $^{-2}$ is the assumed neutron star moment of inertia, P is the rotation period in seconds, and \dot{P} is the period time derivative.

Recent studies indicate that γ -ray emission from pulsars originates at high altitudes in the pulsar magnetosphere near R_{LC} and from the last open field lines (e.g., Abdo et al. 2009b, 2009c, A10). Two classes of high-altitude particle acceleration and

emission model are generally considered: slot-gap accelerators (Muslimov & Harding 2004; Dyks & Rudak 2003) from the polar cap outward along the last open field lines and outer-gap accelerators (Romani & Yadigaroglu 1995; Cheng & Zhang 1998) where the last closed field lines cross the null-charge surface defined by $\Omega \cdot B = 0$ (where Ω is the rotation velocity vector and B is the magnetic field vector). These models result in what we shall term *wide-beam* emission geometries. Low-altitude models for γ -ray emission have been largely ruled out by the generally wide and multi-component γ -ray pulse profiles and the existence of pulsed γ -ray emission at energies greater than a few GeV; at low altitudes such emission is precluded by electron–positron pair absorption. Furthermore, Romani & Watters (2010) find that outer-gap models are statistically preferred over slot-gap models based on fits to LAT pulse profiles for a sample of bright γ -ray pulsars.

Four of the five brightest γ -ray pulsars detected by the Energetic Gamma-Ray Experiment Telescope (EGRET) are also radio pulsars (Thompson 2008). This, together with the similar radio and γ -ray pulse morphology for many young pulsars and MSPs, led Manchester (2005) to suggest that, for these high- \dot{E} pulsars, the radio emission is also emitted in wide beams from the outer magnetosphere near R_{LC} . Within the framework of polar cap emission, Johnston & Weisberg (2006) and Karastergiou & Johnston (2007, hereafter KJ07) proposed that the radio emission from young pulsars is emitted at intermediate altitudes of ~ 1000 km or $\sim 0.2 R_{LC}$ in order to account for wide profiles of young pulsars.

In this Letter, we identify samples of γ -ray-selected and radio-selected pulsars and consider the pulsars in each sample that also emit radio and γ -ray pulses, respectively. We then discuss the implications of our analysis for the radio beams of young pulsars and MSPs.

2. THE γ -RAY AND RADIO PULSAR SAMPLES

We examine two pulsar samples: a γ -ray-selected sample (Sample G) and a radio-selected sample (Sample R). Each

sample consists of a set of pulsars detectable using “blind search” techniques in the relevant band.

2.1. The γ -ray-selected Sample G

A10 presented a catalog of γ -ray pulsars detected using two different techniques. First, 29 γ -ray pulsars were found using timing ephemerides, mostly from contemporaneous radio observations of 218 pulsars with $\dot{E} > 10^{34}$ erg s $^{-1}$ (Smith et al. 2008). Twenty of these were young pulsars and nine were MSPs. The second method involved a blind search for γ -ray pulsars at the positions of ~ 300 unidentified γ -ray sources in the 6 month LAT all-sky image. This search was sensitive to periodicities between 0.15 s and 2 s, and yielded 16 γ -ray pulsar discoveries. A further two pulsars with existing γ -ray rotation ephemerides—Geminga and PSR J1124–5916—were also re-detected. While Geminga is well known for its lack of radio pulses (Kassim & Lazio 1999), PSR J1124–5916 is a radio pulsar for which a γ -ray rotation ephemeris was used because the extremely weak radio emission makes radio timing difficult (Camilo et al. 2002).

Sample G consists of all γ -ray pulsars found in the LAT blind search. We further include those pulsars found using radio ephemerides that have phase-averaged γ -ray fluxes above the blind search sensitivity limit.¹ This is a conservative estimate, as some of the LAT blind search pulsars are in fact below the limit. Our inclusion in Sample G of pulsars where the γ -ray pulses were first detected using radio-based ephemerides is justified because the blind search over the unidentified γ -ray point sources was carried out after the search using radio ephemerides. In total, we included the 18 pulsars found in the γ -ray blind search and 17 pulsars found using radio ephemerides, of which three are MSPs. Searches for radio pulses from 15 of the blind search pulsars by Camilo et al. (2009b) have revealed two detections. We do not discard the three blind search pulsars that have not been searched for radio emission by Camilo et al. (2009b) because their positions have already been searched by the Greenbank low-luminosity pulsar survey (Dewey et al. 1985) with a limiting flux density of 2 mJy at 400 MHz. Therefore, of the 35 γ -ray pulsars in Sample G, 20 are radio pulsars.

2.2. The Radio-selected Sample R

Sample R was formed from the 218 high- \dot{E} radio pulsars observed by the LAT, of which 25 are also γ -ray pulsars. \dot{E} values were corrected for the Shklovskii effect (Shklovskii 1970) using values of \dot{P} , proper motion, and NE2001 distances (Cordes & Lazio 2002) given in the ATNF Pulsar Catalogue² (Manchester et al. 2005); only pulsars with corrected $\dot{E} > 10^{34}$ erg s $^{-1}$ were included in Sample R. We note that the Sample G pulsars have already had the Shklovskii corrections applied (A10).

PSR B1509–58 was not detected in the LAT data, but was found by the *AGILE* γ -ray telescope (Pellizzoni et al. 2009) and so we include it in Sample R. Another pulsar, PSR J0034–0534, was detected using 13 months of *Fermi* LAT data (Abdo et al. 2010a). Thirteen of the high- \dot{E} pulsars were not detected in blind radio surveys, but were found in deep searches of supernova remnants (SNRs; e.g., Camilo et al. 2009a) and EGRET error boxes (e.g., Keith et al. 2008). We discard these 13 pulsars from

Table 1
The γ -ray-selected and Radio-selected Samples

Sample	γ -ray Pulsars	Radio Pulsars
γ -ray-selected (G)	35	20
Radio-selected (R)	17	201

Sample R. This leaves 201 radio-selected pulsars in Sample R, of which 17 are γ -ray pulsars. Table 1 summarizes the two samples.

3. IMPLICATIONS FOR THE RADIO BEAMS

We consider the radio and γ -ray beams as being defined by emissivity functions $F_i(\theta, \phi)$, where the index $i = (r, g)$ indicates, respectively, the radio and γ -ray bands. The co-latitude angle θ is measured from the rotation axis and ϕ is the longitudinal angle. The “latitude coverage,” $C_i(\theta)$, of a beam is defined as the longitude integration of F_i :

$$C_i(\theta) = \int_0^{2\pi} F_i(\theta, \phi) d\phi. \quad (2)$$

The fraction of the celestial sphere swept out by each beam—the “beaming fraction”—is given by

$$f_i = \frac{1}{4\pi} \int_0^\pi \int_0^{2\pi} F_i(\theta, \phi) \sin \theta d\phi d\theta = \frac{1}{4\pi} \int_0^\pi C_i(\theta) \sin \theta d\theta, \quad (3)$$

where the normalization is such that $0 \leq f_i \leq 1$.

To assist with visualization, we approximate the beams by “top-hat” functions where C_i is constant (and non-zero) over some beamwidth Θ_i and zero elsewhere as illustrated in Figure 1. Θ_i is effectively the equivalent width of the beam in latitude. We note that the beaming fraction f_i defined by Equation (3) corresponds to the f_Ω term defined in Watters et al. (2009) for such a top-hat beam. There are four possible configurations for a pulsar emitting in both the radio and γ -ray bands: (1) the radio and γ -ray beams have no overlap; (2) partial overlap of the beams; (3) the radio beam is narrower than the γ -ray beam and entirely overlapping; (4) the γ -ray beam is narrower than the radio beam and entirely overlapping.

If we assume that pulsar rotation axes are randomly oriented in space and that pulses are detectable for observer lines of sight within the beams, then the number of pulsars detectable in band i is proportional to the beaming fraction f_i . Furthermore, provided the beam co-latitude θ_i is not too small, then the number of pulsars detectable in band i is approximately proportional to the beamwidth Θ_i . Importantly, the number detected in both bands is proportional to the width of the overlap region Θ_c .

The issue of pulsar detectability is of course critical. It is well known that the γ -ray luminosity of a pulsar, L_g , is strongly dependent on its spin-down luminosity \dot{E} . While some γ -ray emission models predict $L_g \propto \dot{E}^{1/2}$ (e.g., Harding 1981), we assume an empirical relation $L_g \propto \dot{E}$ after Smith et al. (2008). This implies the commonly used γ -ray detectability metric of the pulsar spin-down flux at the telescope, $\dot{E}d^{-2}$, where d is the pulsar distance. In contrast, the radio luminosity of pulsars is essentially uncorrelated with both \dot{E} and $\dot{E}d^{-2}$. Given that nearly 2000 radio pulsars are known compared to just 46 γ -ray pulsars, it is reasonable to assume that selection effects are much more significant in the γ -ray band and that all γ -ray pulsars emit potentially detectable radio beams.

¹ The LAT blind search sensitivity limit is plotted as a function of the diffuse γ -ray background in Figure 10 of A10.

² <http://www.atnf.csiro.au/research/pulsar/psrcat>

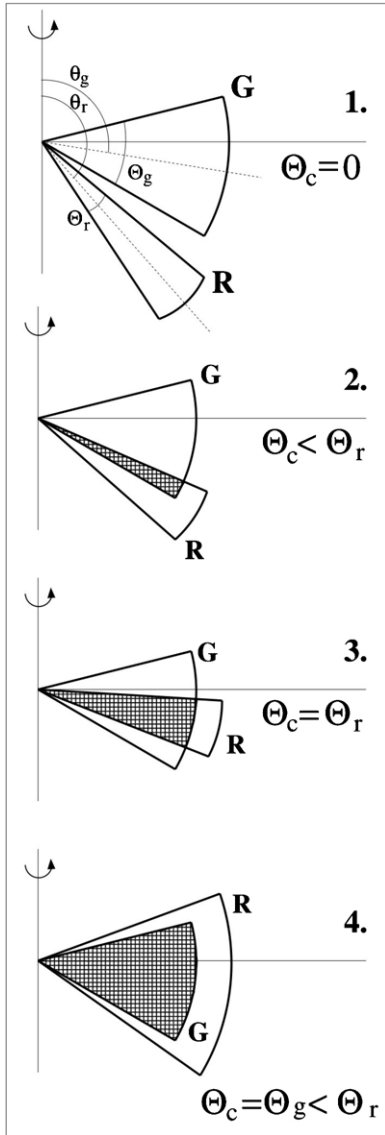


Figure 1. Four possible configurations for the radio and γ -ray beams from a pulsar emitting in both bands. The sectors represent the latitude coverage of the sky as the pulsar rotates for the γ -ray beam (G) and the radio beam (R) where the beams are approximated by “top-hat” functions for clarity. The beams are centered at co-latitudes θ_g and θ_r and have equivalent widths of Θ_g and Θ_r for the γ -ray and radio bands, respectively. The angular zone common to the two beams is shown as a hatched area and has a width Θ_c . If the observer’s line of sight lies within the hatched area, both radio and γ -ray pulses can be observed.

Therefore, for the γ -ray-selected Sample G, the radio fraction, i.e., the ratio of the number of radio detections to the total, $\mathcal{F}_{r,G} = 20/35$, is an estimate of the ratio Θ_c/Θ_g . For configuration 3 in Figure 1, $\Theta_c = \Theta_r$; in all other configurations $\Theta_c < \Theta_r$. Therefore,

$$\frac{f_r}{f_g} \approx \frac{\Theta_r}{\Theta_g} \geq \mathcal{F}_{r,G} \approx 0.571.$$

Since the values of $\mathcal{F}_{r,G}$ follow a binomial distribution, the standard error on this mean value is

$$\sigma = \sqrt{\frac{\mathcal{F}_{r,G}(1 - \mathcal{F}_{r,G})}{n}} = 0.084,$$

where $n = 35$ is the number of pulsars in the sample.

We conclude that, on average, for γ -ray emitting pulsars, the radio beaming fraction is at least half the γ -ray beaming fraction.

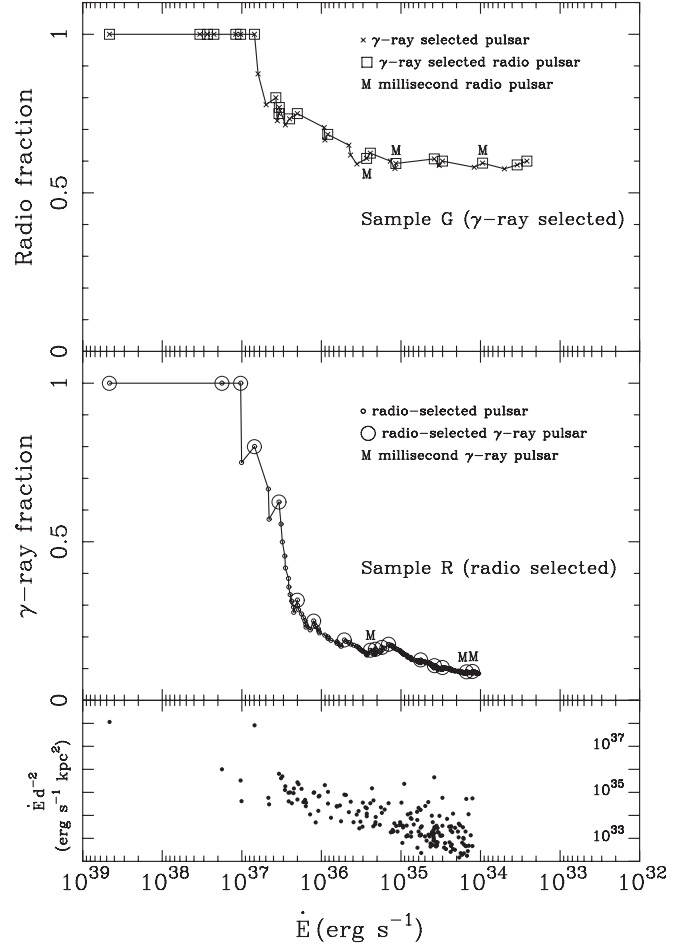


Figure 2. Radio-detected fraction of Sample G, $\mathcal{F}_{r,G}$ (top), and the γ -ray detected fraction of Sample R, $\mathcal{F}_{g,R}$ (middle), computed for each pulsar for the sub-sample of all pulsars having \dot{E} greater than or equal to the value for that pulsar. Note that \dot{E} increases to the left on the plot. The bottom plot shows values of $\dot{E}d^{-2}$ for the Sample R pulsars.

It is possible that some of the current Sample G pulsars will be detected in the radio band, which would increase $\mathcal{F}_{r,G}$. With the continued operation of the LAT, it is certain that the sample of γ -ray-selected pulsars will increase. The value of $\mathcal{F}_{r,G}$ may either increase or decrease, but its uncertainty will decrease. The longitudinal geometries of the beams are not constrained; therefore, the radio and γ -ray emission regions need not be co-located.

To quantify the true latitude coverage of the radio beams with respect to the γ -ray beams, as opposed to just constraining the overlapping area, we compare the radio-selected Sample R with Sample G. First, given that the LAT is only sensitive to pulsars with high values of $\dot{E}d^{-2}$, the γ -ray fraction in Sample R, $\mathcal{F}_{g,R} = 17/201$, is not significant to this discussion. Instead, we order the pulsars in the two samples by decreasing \dot{E} and consider the ratios $\mathcal{F}_{r,G}$ and $\mathcal{F}_{g,R}$ as functions of \dot{E} . For each pulsar, we compute the ratios for the sub-sample of pulsars with \dot{E} greater than or equal to the value for that pulsar.

The results are shown graphically in Figure 2. It is striking that the seven highest- \dot{E} Sample G pulsars are detected in both bands. Within the same \dot{E} range, four of five Sample R pulsars are also γ -ray emitters. It is therefore evident that $\mathcal{F}_{r,G}$ and $\mathcal{F}_{g,R}$ have comparable values in this high- \dot{E} regime. For lower- \dot{E} values and correspondingly lower- $\dot{E}d^{-2}$ values in Sample R (shown in

the bottom panel of Figure 2) $\mathcal{F}_{g,R}$ reduces sharply, but this is entirely due to the $\dot{E}d^{-2}$ selection for γ -ray pulsar detections. However, it is significant that some radio-only pulsars with \dot{E} between 2×10^{36} and 4×10^{36} erg s $^{-1}$ have high values of $\dot{E}d^{-2}$. These pulsars probably have γ -ray beams not directed toward us, suggesting that the radio beams for these high- \dot{E} pulsars are of comparable width to the γ -ray beams and do not always overlap with their latitude coverage. Any of configurations 1, 2, or 4 in Figure 1 could be applicable to this case.

In contrast, under the assumption that radio pulse detectability is independent of \dot{E} , Sample G is free of these selection effects and the drop in $\mathcal{F}_{r,G}$ at $\dot{E} \sim 6 \times 10^{36}$ erg s $^{-1}$ can be interpreted as an evolution in $\mathcal{F}_{r,G}$ from unity for high- \dot{E} γ -ray pulsars to ~ 0.45 for lower- \dot{E} γ -ray pulsars. In general, the highest- \dot{E} pulsars are among the youngest known pulsars, as determined, for example, by SNR associations (Kaspi 2000). Furthermore, pulsars evolve toward lower \dot{E} values. It therefore appears that, for the highest- \dot{E} pulsars, the radio and γ -ray beams have comparable beaming fractions with largely overlapping radio and γ -ray beams. As these pulsars age, the radio beams evolve to overlap about half the latitude coverage of the γ -ray beams. Ultimately, the pulsars become undetectable in the γ -ray band and the radio beams become even narrower as they evolve toward the beaming fractions of ~ 0.1 typical for older pulsars (Tauris & Manchester 1998).

4. MILLISECOND PULSARS

The lower- \dot{E} pulsars in Samples G and R include five MSPs,³ marked with the letter “M” in Figure 2. Despite the differences between normal pulsars and MSPs (e.g., surface magnetic field strength and rotation period) we do not isolate MSPs from our samples for three reasons. First, the three MSPs in Sample G were γ -ray-selected, which imply some commonalities with young γ -ray pulsars. Second, the nine MSPs detected with the LAT by Abdo et al. (2009a, 2010a) have values of $\dot{E}d^{-2}$ that are among those of the top 18 MSPs searched, which give a γ -ray fraction of 0.5, consistent with the lower bound on the radio fraction for the lower- \dot{E} Sample G pulsars (~ 0.45). Third, it has been suggested by Manchester (2005), based on pulse profile and giant pulse characteristics, that MSPs have wide radio beams similar to those of young pulsars. Therefore, in the discussion below on the radio beams of γ -ray emitting pulsars, we treat young pulsars and MSPs together.

5. DISCUSSION

To test the consistency of our results with various empirical polar-cap models for pulsar beaming, we carried out a series of Monte Carlo simulations for the Sample G pulsars to calculate their values of $\mathcal{F}_{r,G}$. Sample R was not considered because of the $\dot{E}d^{-2}$ selection effects discussed above. We assume that the γ -ray beams sweep the entire celestial sphere, that is, we assume $f_g = 1$, consistent with outer-magnetosphere models (e.g., Watters et al. 2009). For each pulsar, 10^5 pairs of α and ζ were randomly chosen from sinusoidal distributions peaked at $\pi/2$ (α is the angle between the rotation and magnetic axes and ζ is the angle between the rotation axis and the line of sight). The radio beam opening half-angle ρ , measured from the magnetic axis, was calculated for each Sample G pulsar and the beaming model. The beaming fraction f_r was then derived

by finding the fraction of points with $|\alpha - \zeta| < \rho$. Since both α and ζ ranged between 0 and $\pi/2$, emission from both magnetic poles was accounted for.

We first consider the model of Rankin (1993) with $\rho = 5.8 P^{-1/2}$ for which we obtain an average beaming fraction of $f_r = 0.46$ for the Sample G pulsars, somewhat lower than our estimate of $f_r/f_g \approx f_r \geq 0.57$. We note that our assumption of $f_g = 1$ is conservative in this context; other similar simulations by Romani & Yadigaroglu (1995) and Harding et al. (2007) using theoretical predictions of f_g based on the outer-gap model find $f_r = 0.3$ and $f_r = 0.1$, respectively, much less than the observed value.

We then considered the model of KJ07 which, for young pulsars, has a radio emission height above the polar caps of ~ 1000 km, along with beaming along the last open field lines. This model results in $f_r = 0.74$, larger than our lower bound of ~ 0.57 . Hence, the KJ07 radio beaming model is consistent with our results. However, the fact that the seven highest- \dot{E} pulsars in Sample G are all observed to emit radio pulses (i.e., have $f_r \sim 1$) is not predicted by the KJ07 model. The product of the radio beaming fractions derived from the KJ07 model of these seven pulsars is 0.22; hence, the KJ07 model has a 22% chance of reproducing this result.

In contrast to the KJ07 model, outer-magnetosphere wide-beam models for the radio emission predict $f_r \sim 1$, consistent with our result that $\mathcal{F}_{r,G}$ is near unity for the highest- \dot{E} pulsars. Besides this, there are other compelling reasons to consider wide-beam models for these high- \dot{E} pulsars (cf. Manchester 2005). First, there is an excess of pulsars with widely separated pulse components or interpulses among young, high- \dot{E} pulsars (e.g., see Figure 1 of Weltevrede & Johnston 2008). Many of these pulsars have pulse profiles very similar to observed γ -ray profiles, with bridge emission between the main components. Furthermore, they have frequency-independent component separations, inconsistent with low-altitude polar-cap models. As was first pointed out by Narayan & Vivekanand (1983), polarization results from many young pulsars indicate beams elongated in the latitude direction, with impact parameters, $|\alpha - \zeta|$, that are much greater than observed pulse widths. This is inconsistent with polar-cap models but fully consistent with outer-magnetosphere models.

We therefore propose that the radio emission region for young high- \dot{E} pulsars is located high in the pulsar magnetosphere, relatively close to the null-charge surface and to the γ -ray-emitting region. The evolution in $\mathcal{F}_{r,G}$ with \dot{E} suggests a slot-gap model for the radio emission and hence that radio profile components represent caustics just as for the γ -ray profiles. This interpretation is also consistent with a transition to low-altitude polar-cap emission for older pulsars.

Not all γ -ray pulsars are seen in the radio and not all high- $\dot{E}d^{-2}$ radio pulsars are seen in γ -rays. This implies that, while the radio and γ -ray emission regions may be in the same general region of the pulsar magnetosphere, they are not necessarily co-located. In the case of PSR J0034–0534, however, the radio and γ -ray profiles are almost identical and aligned in phase, suggesting that the emission regions are in fact co-located (Abdo et al. 2010a).

6. CONCLUSIONS

By comparing the relative occurrence of radio and γ -ray pulsed emission in γ -ray-selected and radio-selected samples, we find that the radio beaming fraction f_r appears to be close to unity for the highest- \dot{E} pulsars, decreasing to ~ 0.5 for

³ While Abdo et al. (2009a) report the γ -ray detections of eight MSPs, three did not fall within either of the cutoffs applied to Sample G or R.

lower- \dot{E} γ -ray pulsars. Our estimated lower bound f_r , averaged over all γ -ray-selected pulsars, of ~ 0.57 is inconsistent with low-altitude polar-cap models for the radio emission but is consistent with KJ07's intermediate-altitude polar-cap model. However, the KJ07 model is marginally inconsistent with our result that $f_r \sim 1$ for the highest- \dot{E} pulsars. With other evidence from radio and γ -ray pulse morphologies, these results suggest that, for high- \dot{E} pulsars, the radio emission originates in wide beams from the vicinity of the null-charge surface, possibly with a slot-gap configuration, with profile components representing caustics in the emission pattern.

Improved statistics from further γ -ray and radio discoveries will help to test our conclusions. Modeling of radio pulse profiles with slot-gap or other wide-beam emission geometries for both young high- \dot{E} pulsars and MSPs would be valuable. If outer-magnetosphere wide-beam radio emission from these pulsars is confirmed, it would have profound implications for pulsar population and evolution studies as well as our understanding of radio and γ -ray pulse emission mechanisms.

We thank Simon Johnston and Kyle Watters for valuable discussions and an anonymous referee for constructive suggestions. G.H. is supported by an Australian Research Council QEII Fellowship (project DP0878388).

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